

# Electric-field-induced ferromagnetic resonance in magnetic tunnel junctions

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## Motivation

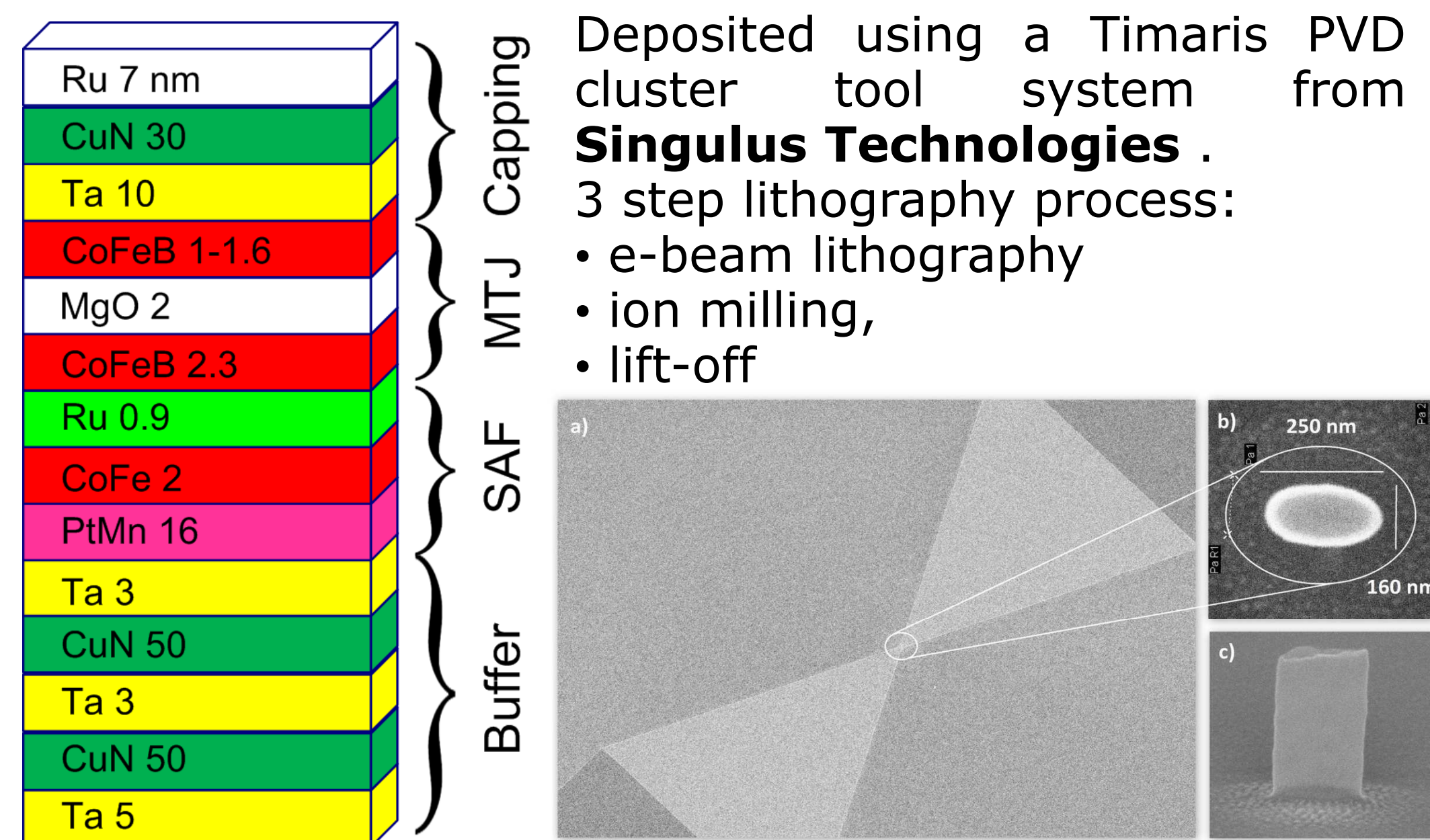
Steady-state magnetization precessions in magnetic tunnel junctions (MTJ) that are induced by microwaves result in a DC response when the precessing and stimulating signals are in resonance [1]. Depending on the MTJ multilayer structure, the magnetization precession can originate either from the spin-transfer-torque (STT) effect [2] or from electric-field-induced (E-field) changes to the magnetic anisotropy of MTJ free layer [3].

## Aims

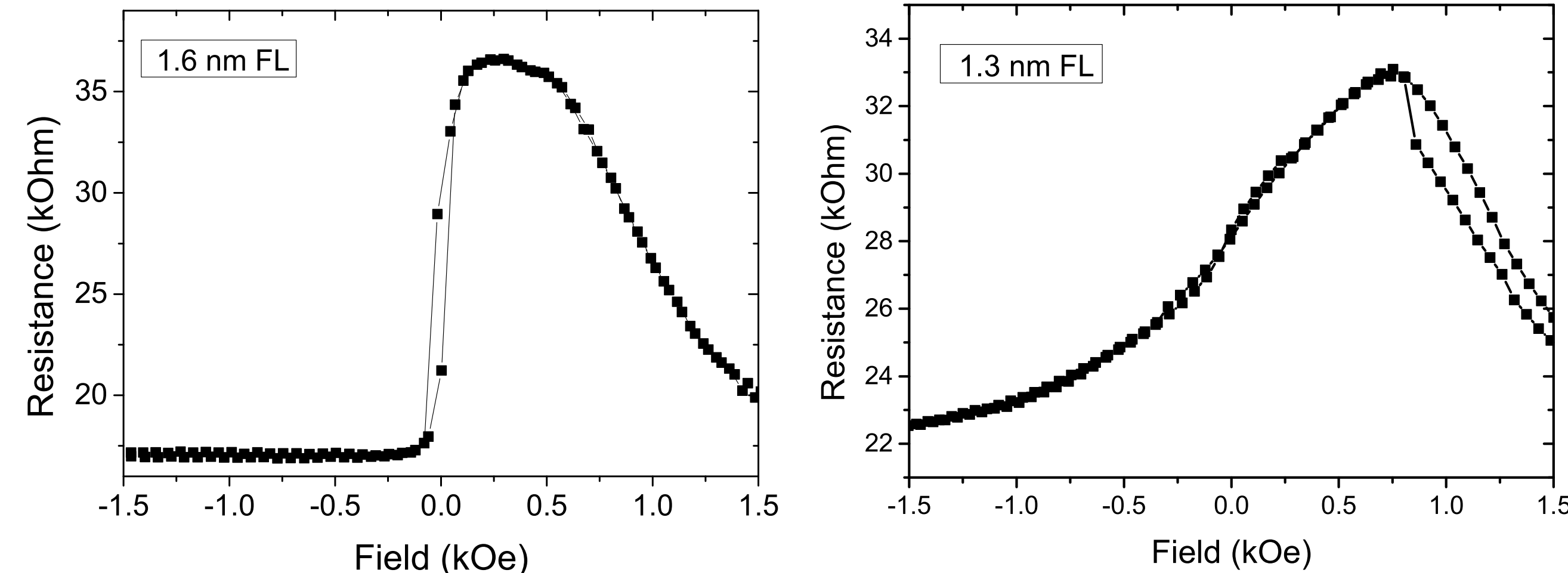
To investigate ferromagnetic resonance (FMR) in MTJ that can be used in microwave oscillators, detectors, magnetic field sensors. Distinguish between: **E-field-FMR** and **STT-FMR**.

## Sample description

### A1. MTJ nanopillar

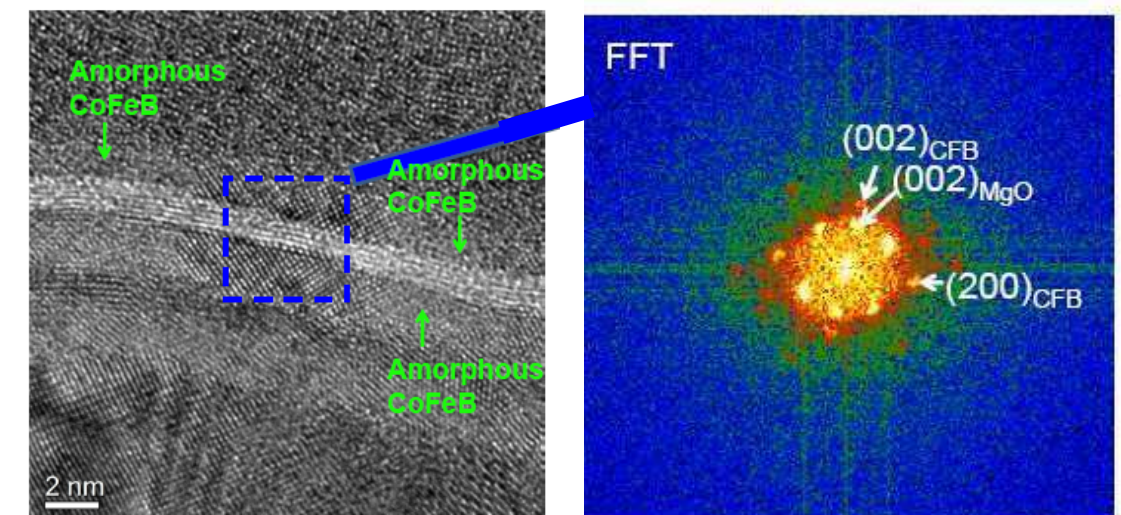


### A2. MTJ optimization



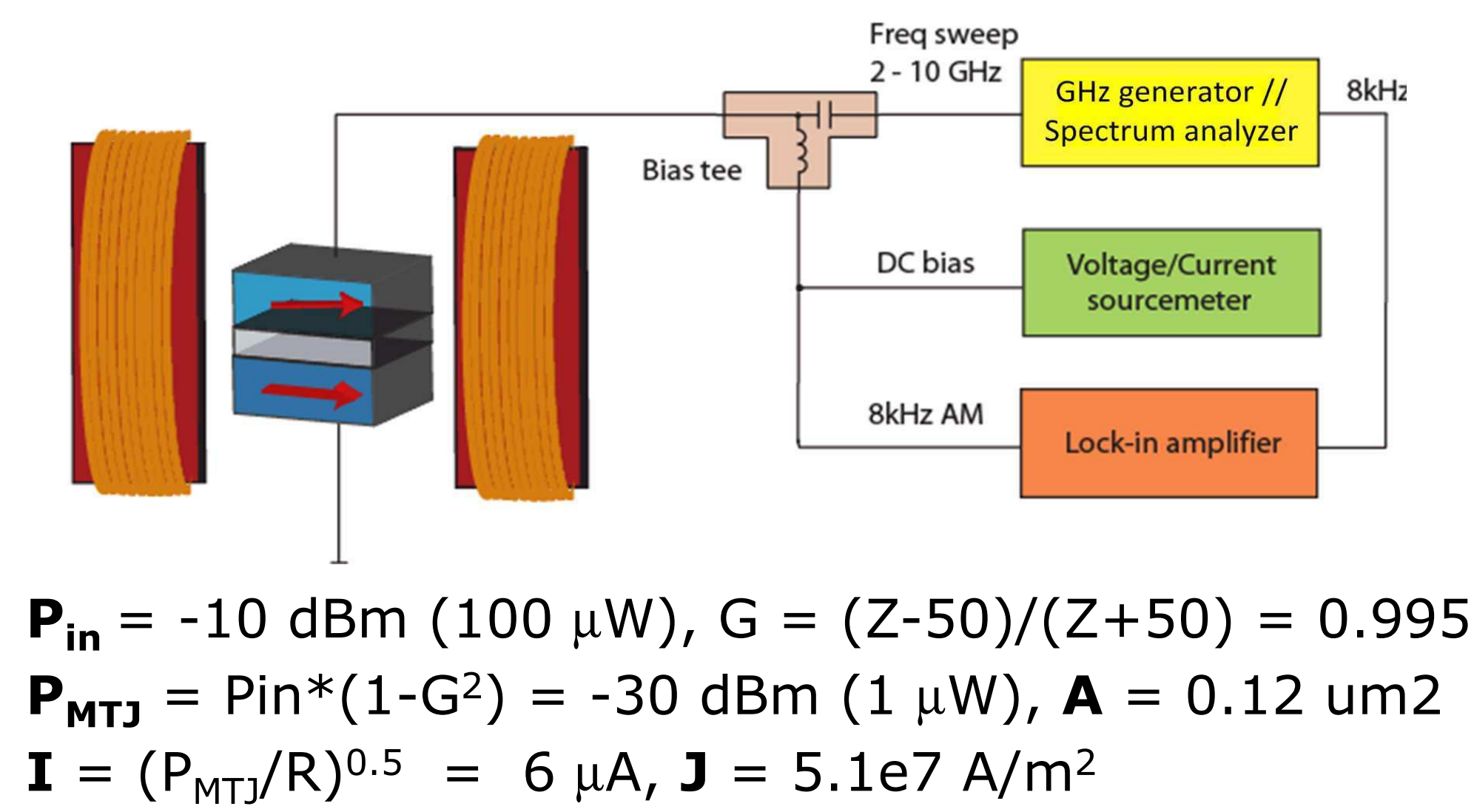
**In-plane** anisotropy for 1.6 nm thick free layer (FL) and **perpendicular** anisotropy for 1.3 nm thick FL. The measured **TMR** ratio is between 70 and 100 %.

2 nm MgO tunnel barrier results in an **RA** product of **2 kΩμm<sup>2</sup>**, two MTJ size of 120 x 230 nm and 280 x 530 nm. Annealed at 250°C with in-plane magnetic field (4 kOe) in order to set the exchange bias direction.

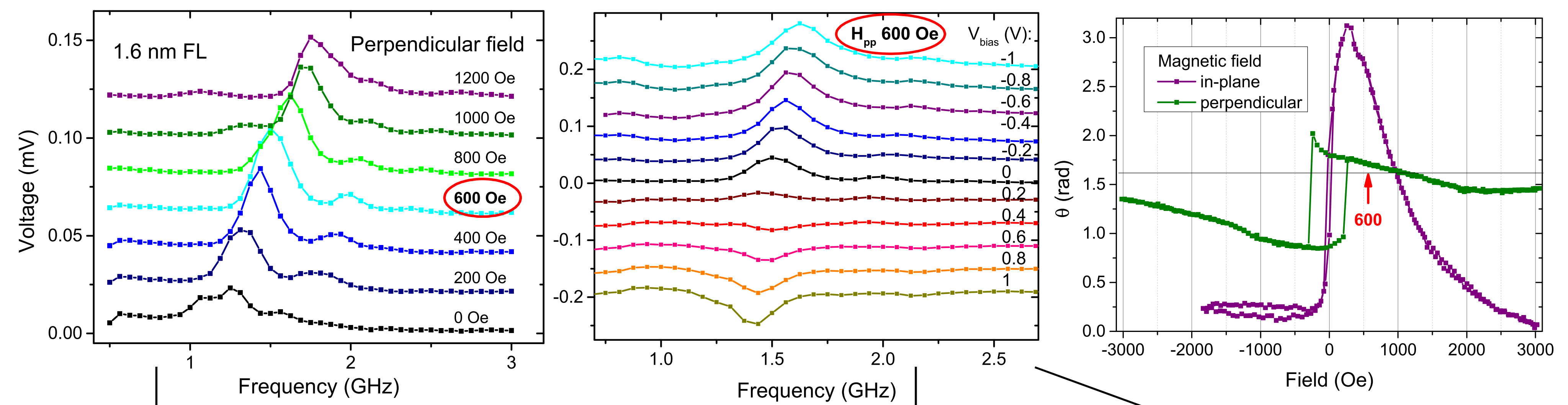


## Experiment and results

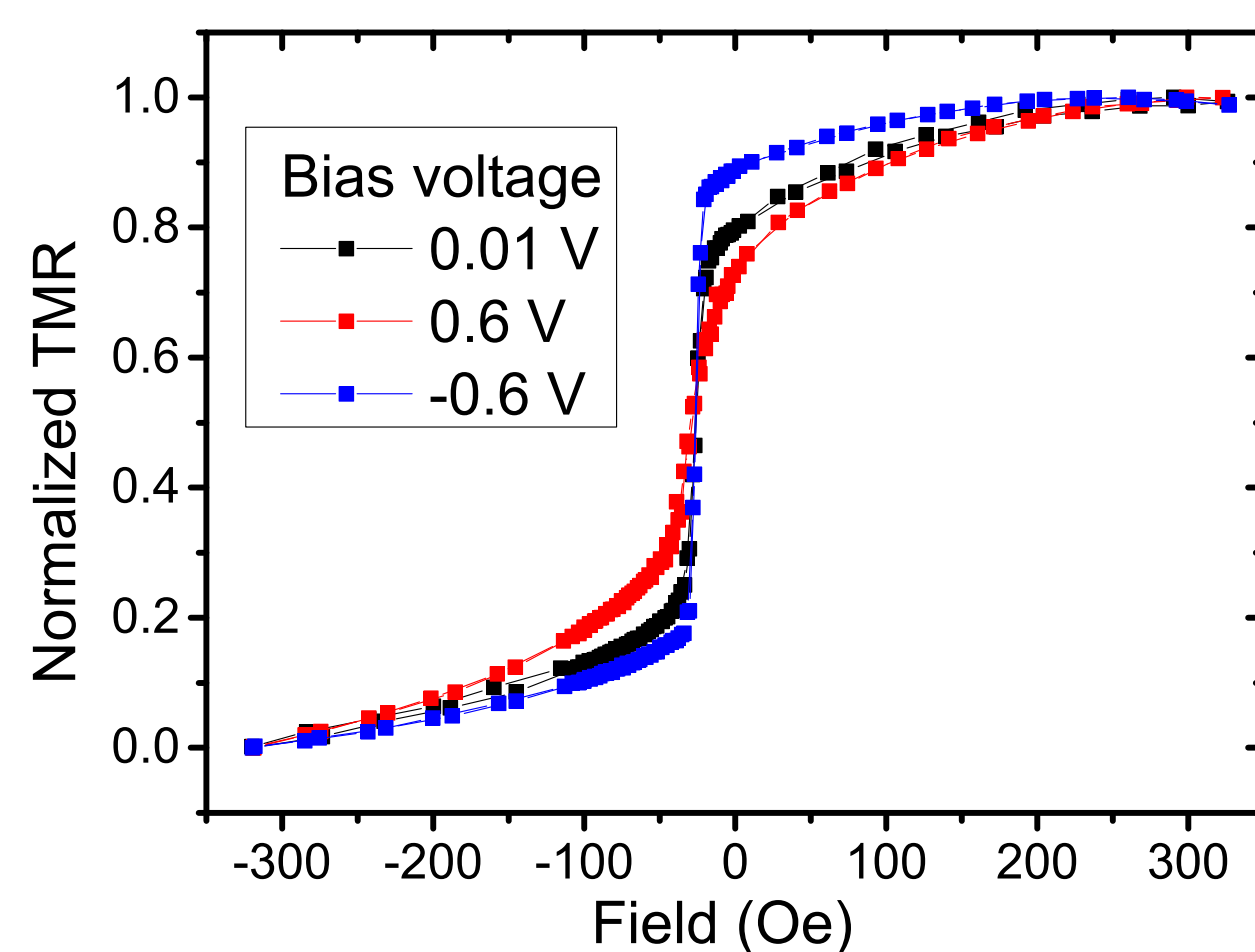
### B1. Measurement setup



### B3. FMR measurements and calculations



### B2. TMR for diff. voltage



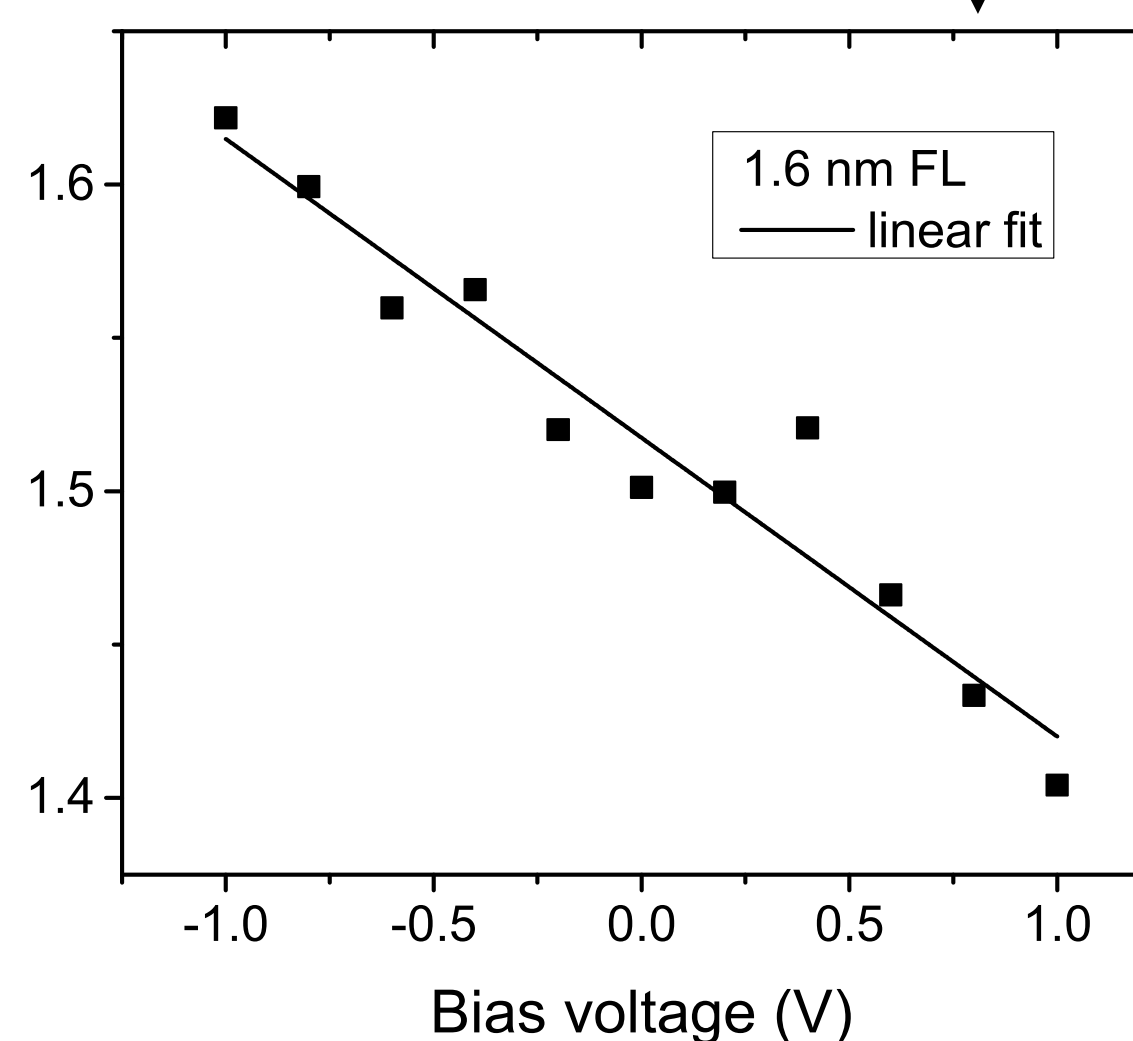
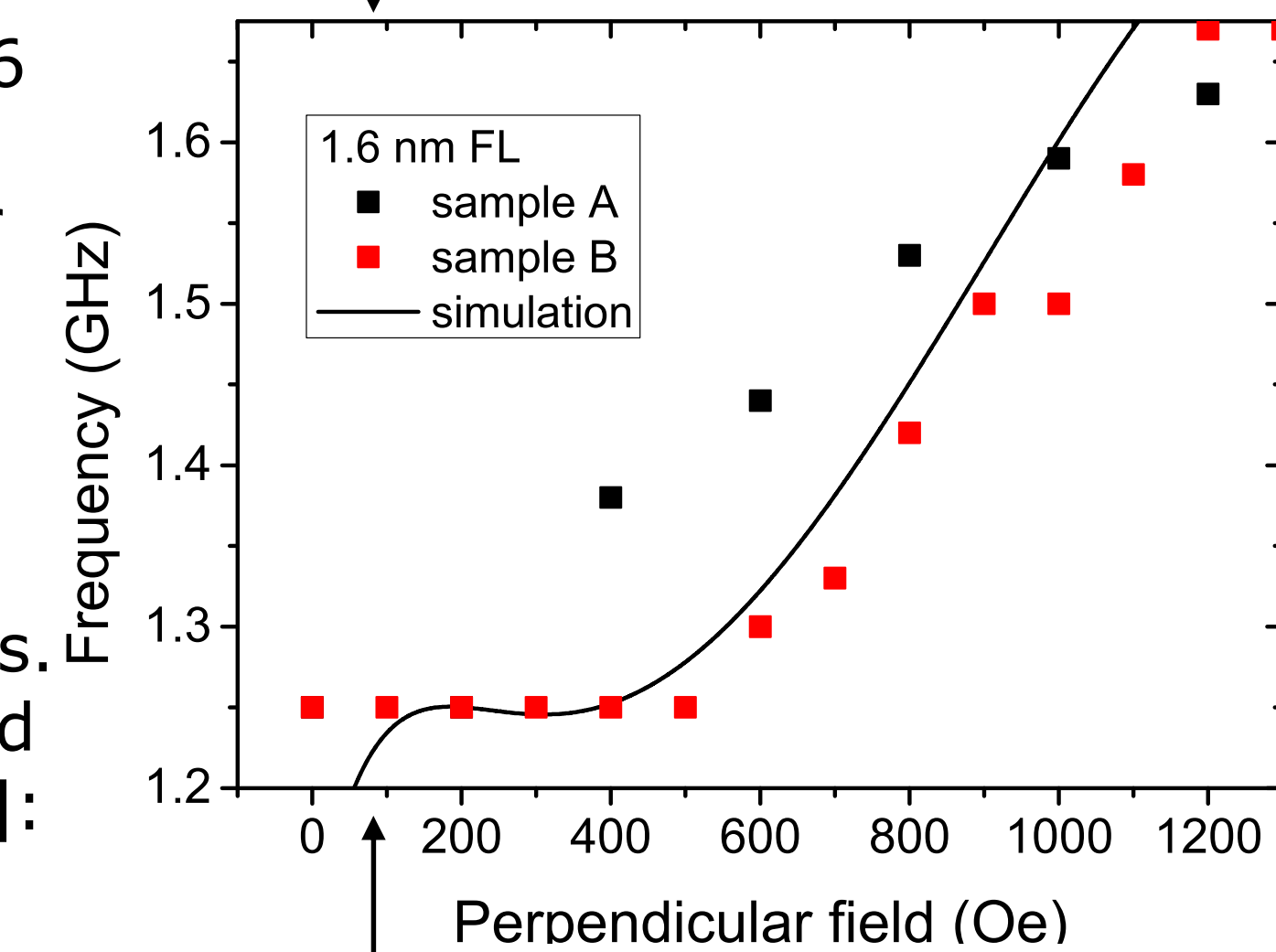
- For MTJs with 1.6 nm FL strong FMR signal detected for perpendicular field
- FMR frequency shifts with bias voltage
- FMR frequency vs. field curve modeled using equation [5]:

$$f = \frac{\gamma_e}{2\pi} \left[ \frac{1}{\Gamma \sin \theta_s} [2H_s \cos 2\theta_s - H \cos(\theta_H - \theta_s) + M_s \cos 2\theta_s (N_{dx} - N_{dy})] \times [-H \sin \theta_H \sin \theta_s + M_s \sin^2 \theta_s (N_{dy} - N_{dx})] \right]^{0.5}$$

### Summary

- MTJ with thick MgO barrier and different FL fabricated and patterned to nm-scale devices
- Magnetic anisotropy shifts from **in-plane** to **perpendicular** at 1.4 nm thick CoFeB layer
- Perpendicular magnetic anisotropy is increased by positive voltage and decreased by negative bias voltage
- FMR signal induced by **STT** is modulated by static bias voltage, in-plane torque is measured in a high bias voltage range

### B4. FMR frequency vs. magnetic field and bias voltage, in-plane STT derivation



Comparison between E-field FMR and STT-FMR:

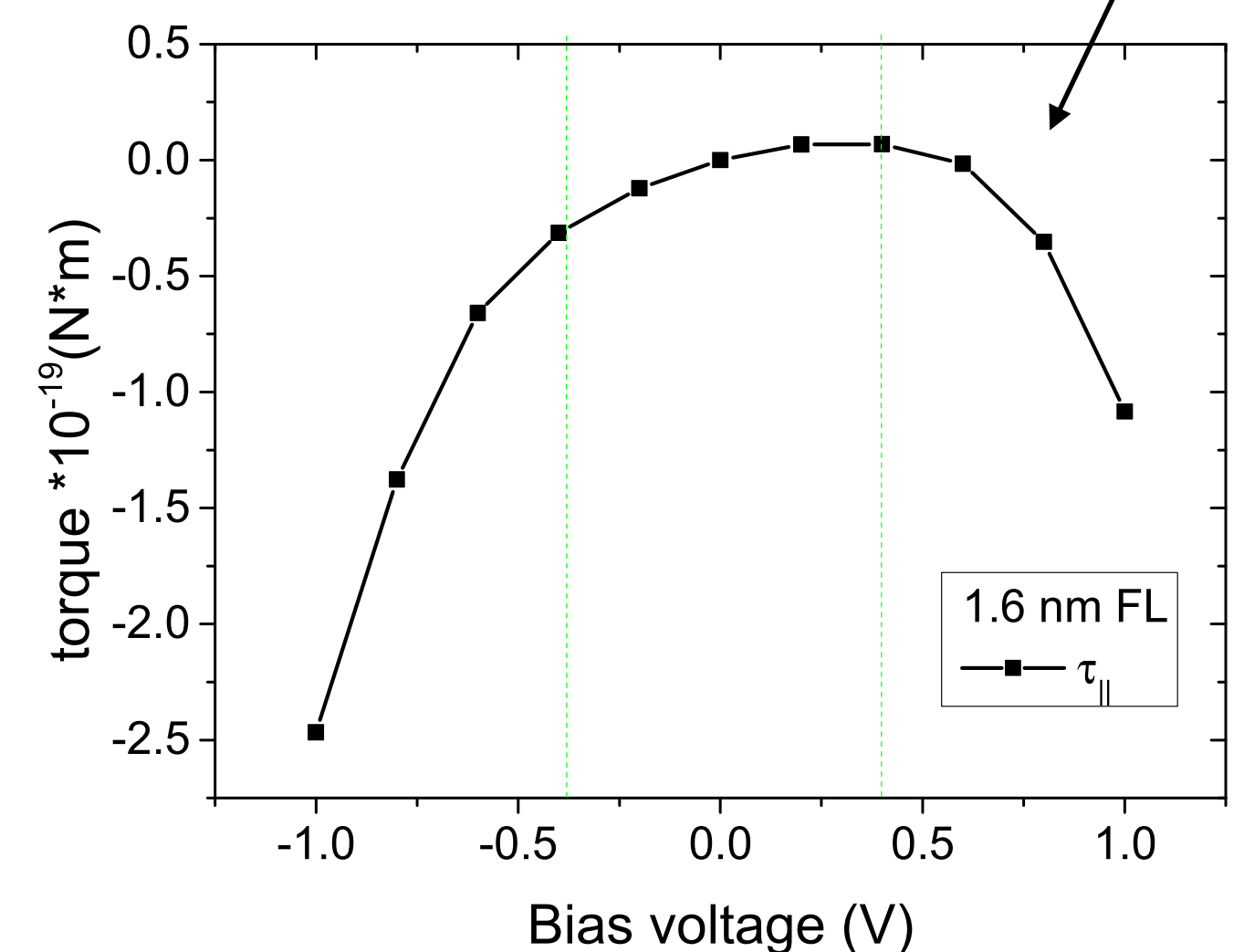
- Different  $V_{mix}$  shape (symmetric vs. asymmetric Lorentzian curve)
- Different angle behaviour (maximum vs. minimum at  $\theta = 90$  degree)
- Different current density supplied (high vs. small)

In-plane STT derived from  $V_{mix}$  signal:

$$V_{mix} = \frac{1}{4} \frac{\partial^2 V}{\partial I^2} I_{RF}^2 + \frac{1}{2} \frac{\partial^2 V}{\partial I \partial \theta} \frac{\hbar \gamma \sin \theta}{4eM_s Vol \sigma} I_{RF}^2 [\xi_{\parallel} S(\omega) - \xi_{\perp} \Omega_{\perp} A(\omega)]$$

Integrated in-plane torque:

$$\xi_{\parallel} = \frac{2e}{\hbar} \sin \theta \frac{dV}{dI} \frac{d\tau_{\parallel}}{dV}$$



**In-plane torque** was measured in a **high** bias voltage range (+/- 1V) revealed parabolic shape [6].

## References and acknowledgments

- [1] – A. A. Tulapurkar et al. Nature 35, 339 (2005) [2] – W. Skowronski et al. PRB 87, 094419 (2013) [3] – T. Nozaki et al. Nature Phys. 8, 491 (2012)  
[4] – W. Skowronski et al. APL 101, 192401 (2012) [5] – J. Zhu et al. PRL 108, 197203 (2012) [6] – M. Wilczyński et al. PRB 77, 054434 (2008)

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